The ρ_{ν_i} increase. Therefore the ρ_e decrease. Due to this decrease, at a given time t temperature T for the models is smaller than T without neutrino heating. So Figures (3) and (4) are graphed vs. the time rather than T. T for the case without heating is shown on the top border.

Figure (3) shows $X_{^4He}$ (the differences between the models is imperceptable on the graph's scale) and $\Delta X_{^4He}(OB, MB) = X_{^4He}(OB, MB) - X_{^4He}$. For most of the run, $X_{^4He}(OB)$ is slightly smaller than $X_{^4He}(MB)$ (too small to see on the graph) perhaps because of the One Bin model's smaller ρ_ν's and slightly larger T. But just after Nucleosynthesis at around $T = 7$ GK $X_{^4He}(OB)$ surpasses $X_{^4He}(MB)$ leading to a larger final result.

$$\begin{aligned}
X_{^4He} &= 0.239272 \\
X_{^4He}(OB) &= 0.239736 \\
X_{^4He}(MB) &= 0.239432
\end{aligned}$$

$\Delta X_{^4He}(OB)$ seems to flatten out earlier than $\Delta X_{^4He}(MB)$. Perhaps the higher energy bins in the MultiBin model convert more neutrons to protons than in the One Bin model.

Figure (4) shows $\lambda(n \to p)$ and $\lambda(p \to n)$ (Again, the differences between models are imperceptable.) and $\Delta\lambda(n \to p, OB, MB) = \lambda(n \to p, OB, MB) - \lambda(n \to p)$, $\Delta\lambda(p \to n, OB, MB) = \lambda(p \to n, OB, MB) - \lambda(p \to n)$. The timestep size makes quantities difficult to calculate at the highest temperatures. So $\Delta\lambda$ is initially negative in the One Bin model while it varies considerably in the MultiBin model. But eventually both models increase their λ's. Moreso for the One Bin model at first. But the increases of $\lambda(n \to p, OB)$ and $\lambda(p \to n, OB)$ are more equal than for the MultiBin model. Eventually $\lambda(n \to p, OB)$ falls below $\lambda(n \to p, MB)$, and is less at the time when Nucleosynthesis occurs. A lower rate means fewer neutrons have decayed. But this difference could be too small to fully account for the higher $X^{^4He}$ in the One Bin model.