$$\begin{aligned}
\rho_{\nu_e}(OB) &= 1.0076\rho_{\nu_0} \\
\rho_{\nu_e}(MB) &= 1.0093\rho_{\nu_0} \\
\rho_{\nu_\mu}(OB) &= 1.0035\rho_{\nu_0} \\
\rho_{\nu_\mu}(MB) &= 1.0038\rho_{\nu_0}
\end{aligned}$$

Equation (6) is equivalent to multiplying the $f_\nu(p)$ by

$$\begin{aligned}
f_{\nu_i}(p) &\rightarrow f_\nu(p)\int_{-\infty}^{\alpha} d\alpha' \left(R'\frac{T}{T_N}\right)^3 W_\nu(\alpha';\alpha) \\
f_{\nu_i}(p) &\rightarrow f_{\nu_i}(p)\frac{\rho_{\nu_i}}{\rho_{\nu_0}}
\end{aligned}$$

Figure (2) shows the increase of $f_{\nu_e}(p_i)$ from the Fermi-Dirac value $f_{FD}(p_i)$ for $p_i = 1$ to 25, with the One Bin increase marked on the right. Only a few of the bins have a change equal to or less than the One Bin change. The higher energy bins have a much larger energy transfer from the $e^\pm$ pairs, raising ρ_{ν_e} over the One Bin result. But note that the f/f_{FD} ratios reach a final value, once the rates involving that $f_{\nu_e}(p_i)$ are no longer effective. At T slightly over 1 GK all the bins large enough to affect the value of ρ_{ν_e} have reached their final values, and so ρ_{ν_e} flattens out.

For higher energy bins the reactions stay effective longer. But near $T = 1$ GK the higher $f_{\nu_e}(p_i)$'s have a sudden decrease. Possibly, the e^-e^+ annihilation rate $e^-e^+ \rightarrow \nu_e\bar{\nu}_e$ has ceased to be effective then. But rates that deplete the higher energy neutrinos into lower energy neutrinos are still working. So $f_{\nu_e}(p_i)$ decreases until these destruction rates themselves are no longer effective. The situation is analogous to the deuterium mass fraction X_d. I think, though that bins higher than 25 would replenish the decreasing bins, reducing the decrease. So this decrease might be an indicator of the BBN code's limitations.

f_{ν_μ} has similar behavior, but much less increase. Also, its highest bins don't get nearly high enough to undergo depletion.